

Simulating NNLO QCD corrections for processes with giant K factors

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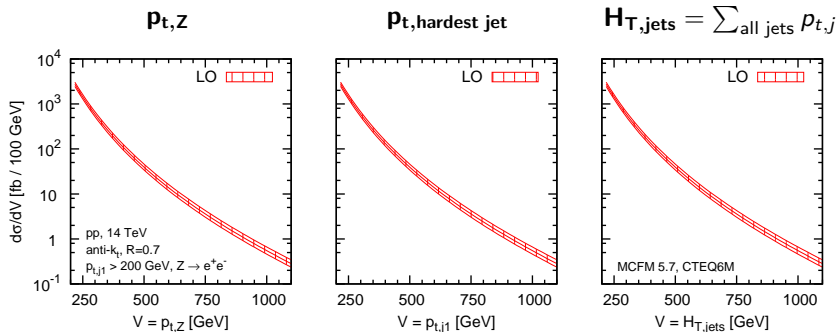
in collaboration with Gavin Salam and Mathieu Rubin¹

HP²3rd, Florence, 14-17 September 2010

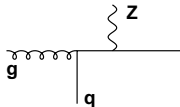
¹M.Rubin, G.P.Salam and SS, arXiv:1006.2144 [hep-ph]

The problem of giant K factors

► Z+j at the LHC

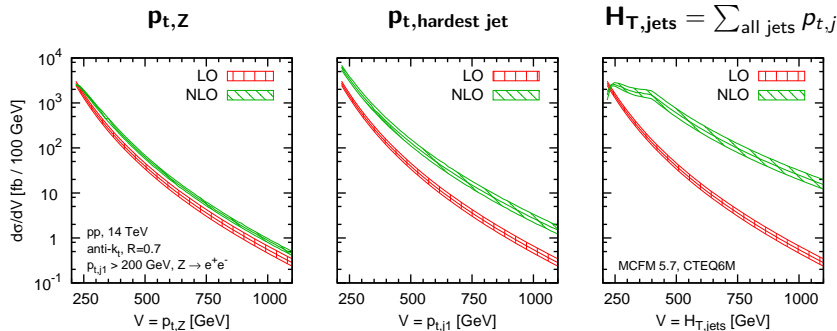


LO:



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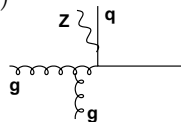
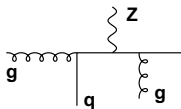
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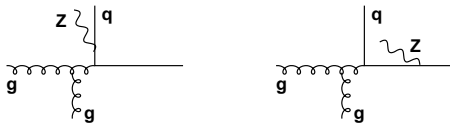
$$\mathcal{O}(\alpha_{\text{EW}}\alpha_s^2 \ln^2 p_{t,j1}/M_Z)$$

NLO:



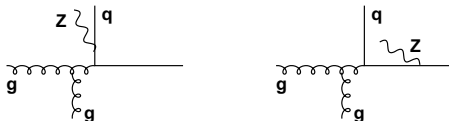
What do we have and what is missing?

- ▶ The large K factor for the Z +jet comes from the new “dijet type” topologies that appear at NLO



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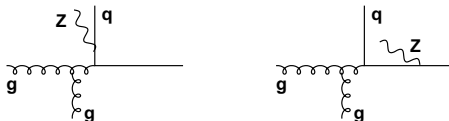
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- ▶ this raises doubts about the accuracy of these predictions
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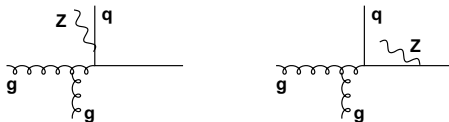


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▶ 2-loop part

- ▶ we need it to cancel IR and collinear divergences from Z+2j at NLO result
- ▶ it will have the topology of Z+j at LO so it will not contribute much to the cross sections with giant K-factor

The basic idea

How to cancel the infrared and collinear singularities?

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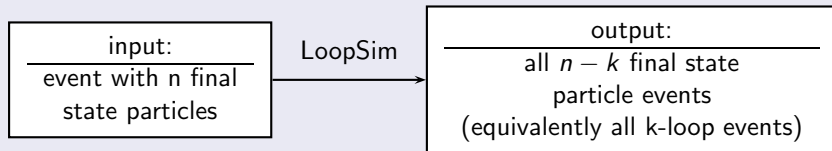
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LoopSim procedure

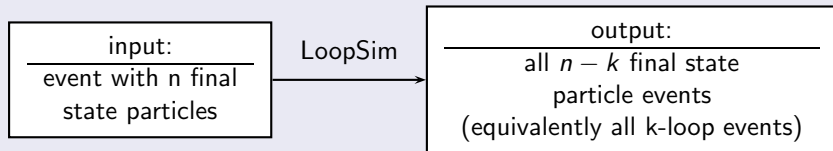


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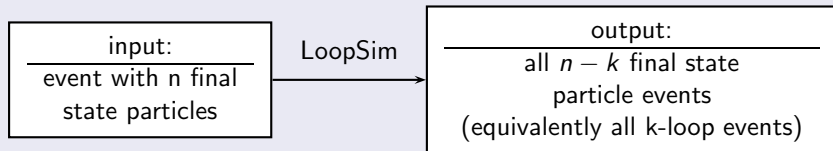
- ▶ notation:
 - $\bar{n}\text{LO}$ – simulated 1-loop
 - $\bar{n}\bar{n}\text{LO}$ – simulated 2-loop and simulated 1-loop
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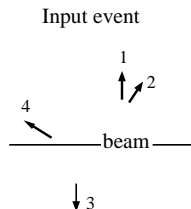
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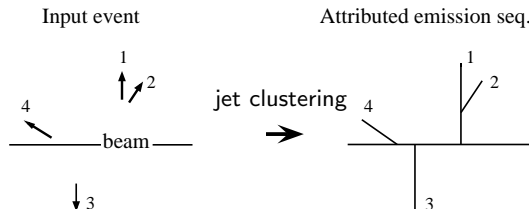
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- ▶ this will work very well for the processes with large K factors e.g.

$$\sigma_{\bar{n}\text{NLO}} = \sigma_{\text{NNLO}} \left(1 + \mathcal{O} \left(\frac{\alpha_s^2}{K_{\text{NNLO}}} \right) \right), \quad K_{\text{NNLO}} \gtrsim K_{\text{NLO}} \gg 1$$

The LoopSim method: \bar{n} LO, $\bar{n}\bar{n}$ LO etc.

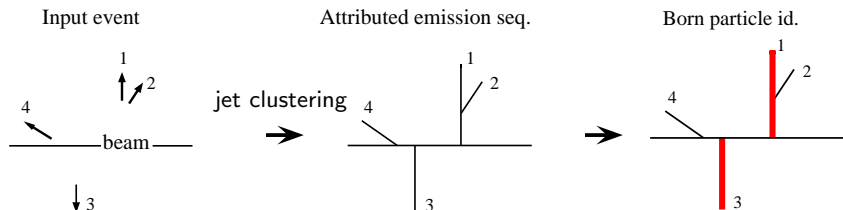


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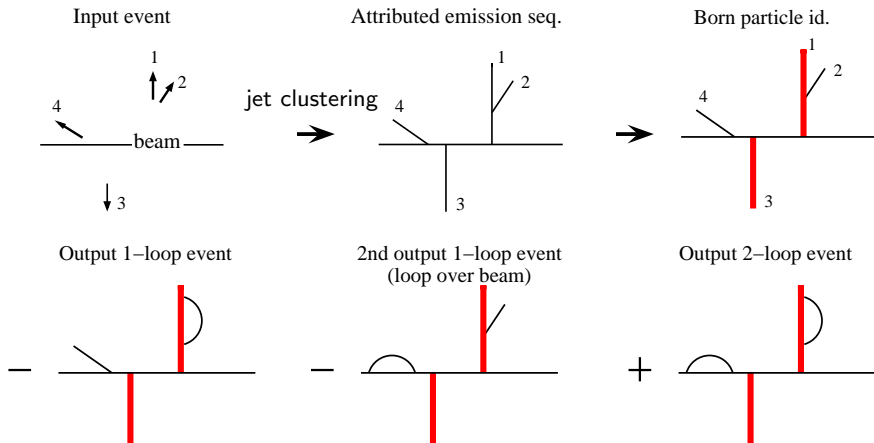
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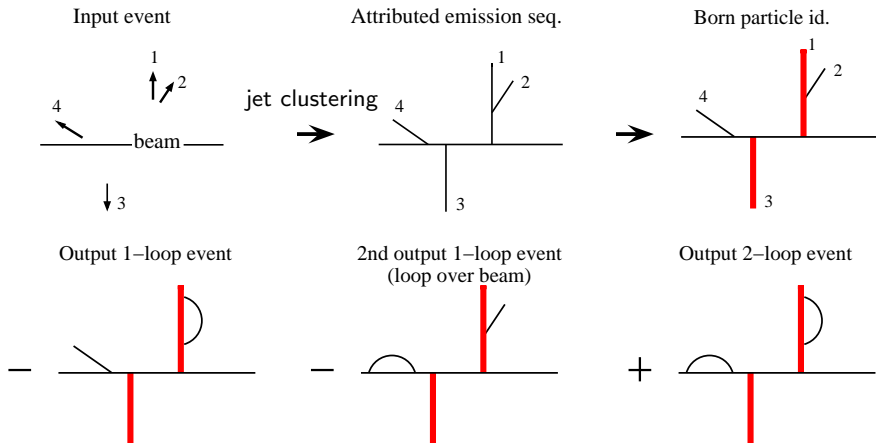
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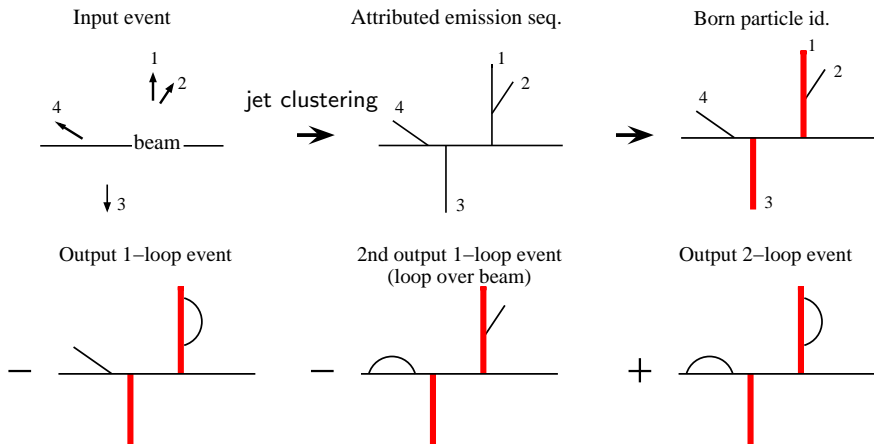
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- ▶ beware: the loops above are just a shortcut notation!

Including exact loops

- $E_{n,l}$ – input event with n final state particles and l loops
- U_l^b – operator producing event with b Born particles and l loops
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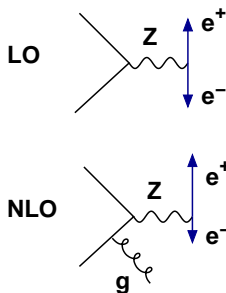
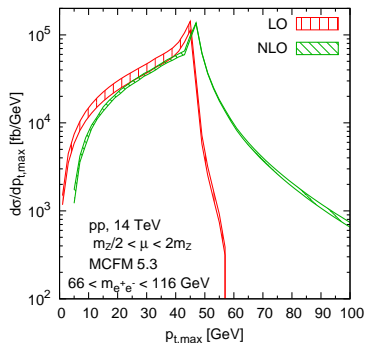
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- ▶ inclusion of exact loops helps reducing scale uncertainties
 - ▶ straightforward generalization to arbitrary number of exact loops

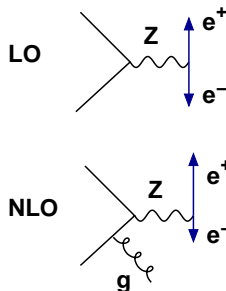
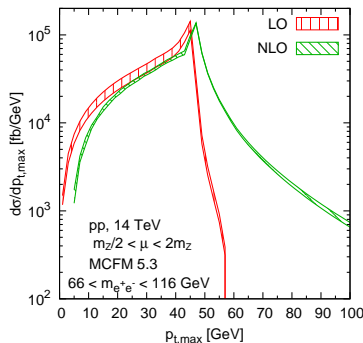
Validation

Drell-Yan at NNLO: spectrum of harder lepton



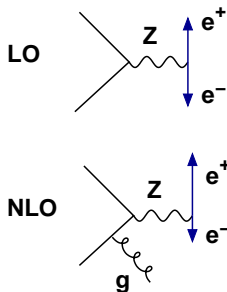
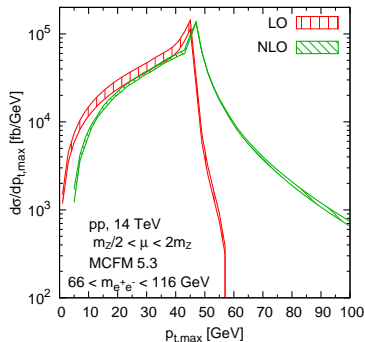
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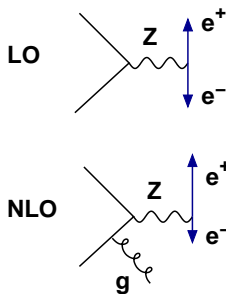
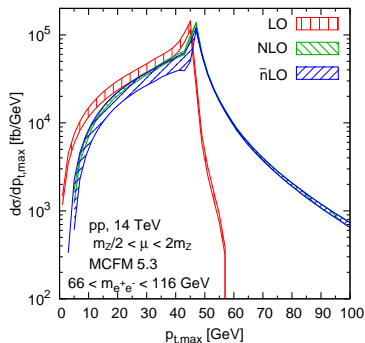
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- ▶ three regions of $p_{t,max}$: $\lesssim \frac{1}{2} M_Z$ $[\frac{1}{2} M_Z, 58 \text{ GeV}]$ $> 58 \text{ GeV}$

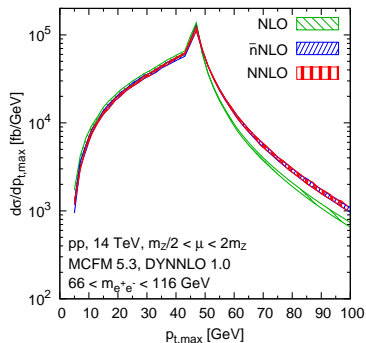
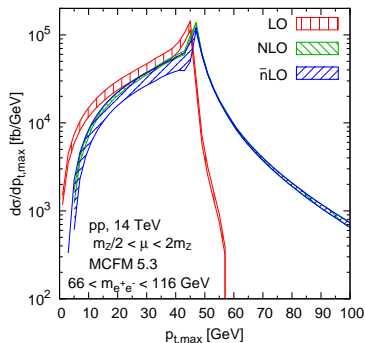
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|--------------------------------|-------------------------------|-------------------------------------|-----------------------|
| $\bar{n}LO$ vs NLO | very good (not guaranteed) | excellent (expected) | perfect (expected) |

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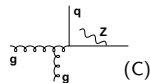
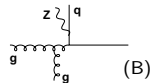
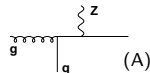
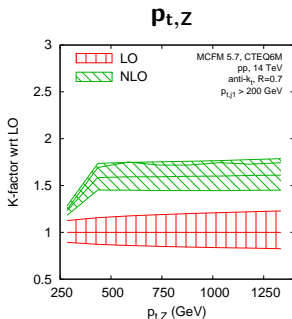


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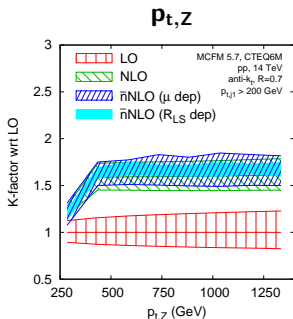
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\bar{n} NLO predictions for LHC

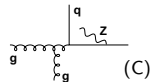
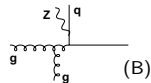
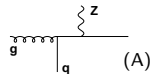
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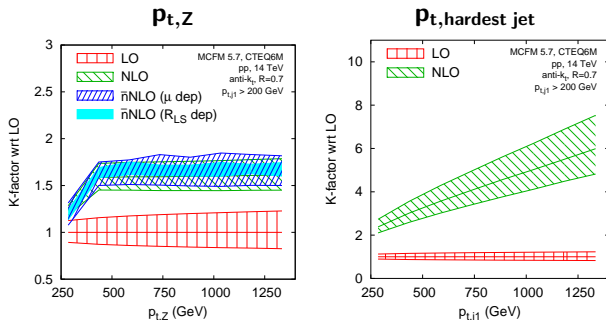
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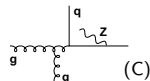
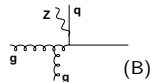
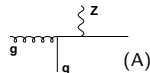
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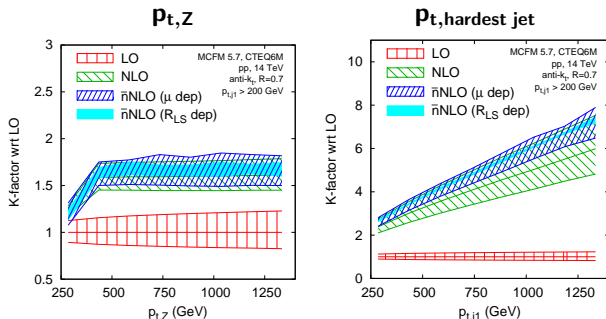
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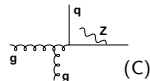
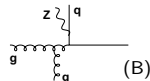
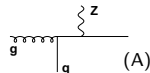
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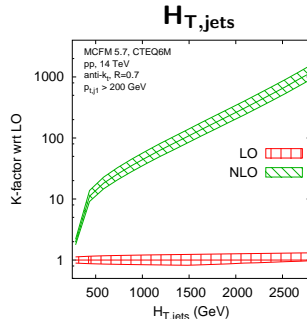
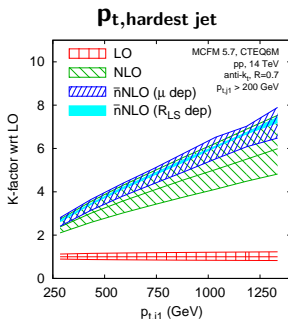
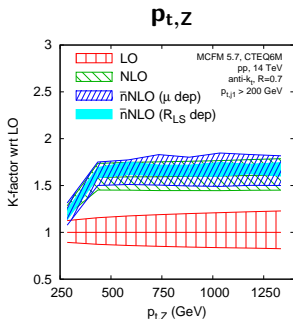
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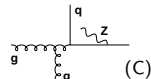
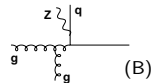
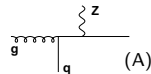
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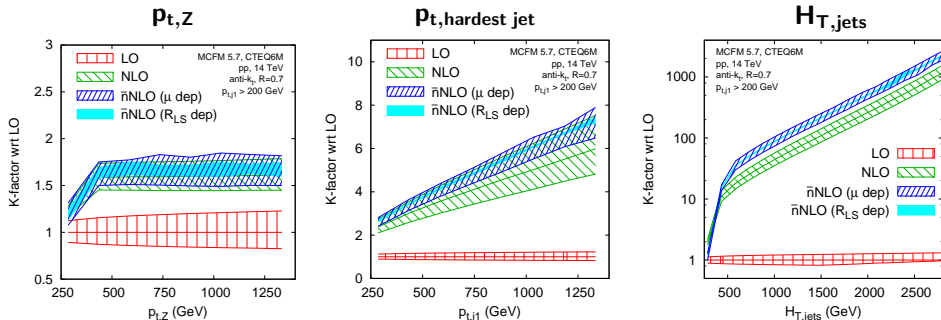
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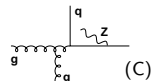
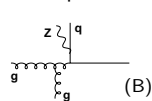
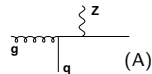
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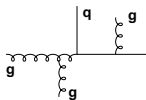
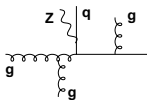


- ▶ $p_{t,Z}$: no correction; topology (A) dominant at high $p_{t,Z}$ (extra loops w.r.t. NLO do not change much)
- ▶ $p_{t,j}$: small correction; $\bar{n}\text{NLO}$ is like NLO for the dominant (B) and (C) configurations and it behaves like healthy NLO
- ▶ $H_{T,jets}$: significant correction; K factor ~ 2 ; given that it is more like going from LO to NLO this may happen sometimes, especially for nontrivial observables like H_T ; can we understand it here?



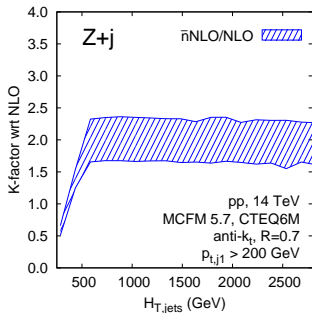
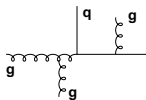
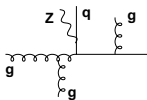
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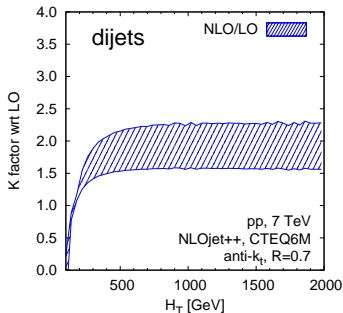
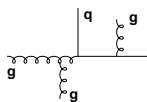
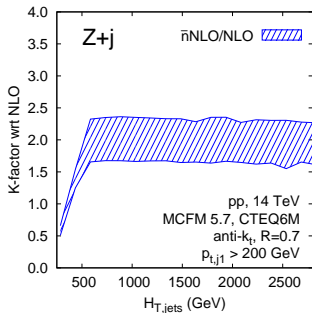
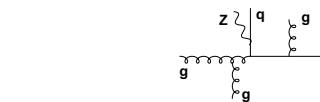
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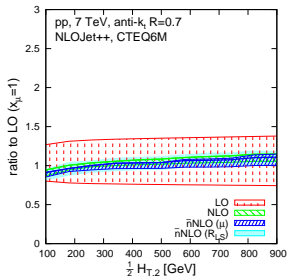


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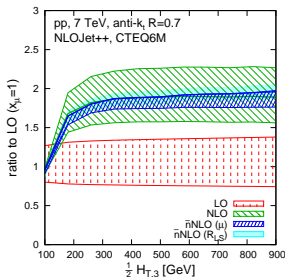
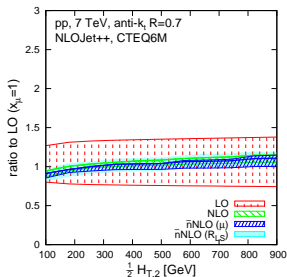
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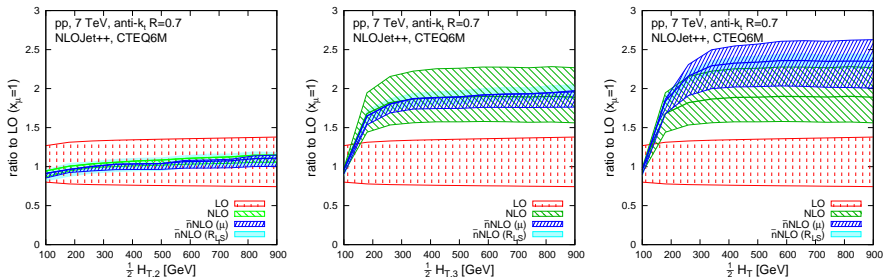
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- ▶ **$H_{T,3}$ converges, significant reduction of scale uncertainty:** the observable comes under control at \bar{n} NLO
- ▶ **H_T does not converge:** again caused by the initial state radiation, this time a second emission which shifts the distribution of H_T to higher values and causes no effect for the $H_{T,3}$ distribution

Summary

- ▶ several cases of observables with giant NLO K factor exist
- ▶ those large corrections arise due to appearance of new topologies at NLO
- ▶ we developed a method, called *LoopSim*, which allows one to obtain approximate NNLO corrections for such processes
- ▶ the method is based on unitarity and makes use of combining NLO results for different multiplicities
- ▶ we gave arguments why the method should produce meaningful results and we validated it against NNLO Drell-Yan and also NLO Z+j and NLO dijets
- ▶ we computed approximated NNLO corrections to Z+j and dijets at the LHC finding, depending on observable, either indication of convergence of the perturbative series or further corrections
- ▶ the latter has been understood and attributed to the initial state radiation

Outlook

- ▶ processes with W , multibosons, heavy quarks, ...

BACKUP SLIDES

The LoopSim method: some more details

For a given input E_n event with n final state particles the weights of all diagrams generated by LoopSim sum up to zero (unitarity)

$$\sum_{\text{all diagrams}} w_n = \sum_{\ell=0}^v (-1)^\ell \binom{v}{\ell} = 0, \quad \ell - \text{number of loops, } v - \text{maximal } \ell$$

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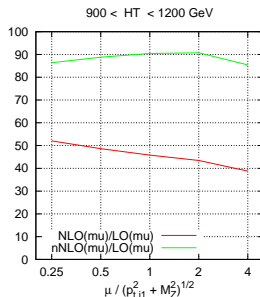
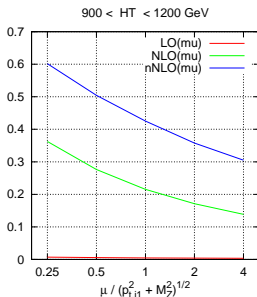
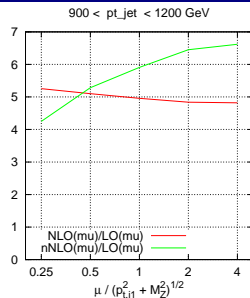
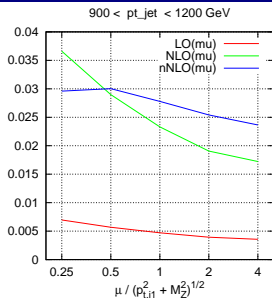
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The principle of the method is simple. There is, however, a number of issues that need to be addressed to fully specify the procedure and make it usable:

- ▶ infrared and collinear safety
- ▶ conservation of four-momentum
- ▶ choice of jet definition (algorithm, value of R)
- ▶ treatment of flavour (e.g. for processes with vector bosons)
 - ▶ Z boson can be emitted only from quarks and never itself emits
- ▶ extension to input events with exact loops

Scale dependence: $Z + \text{jet}$



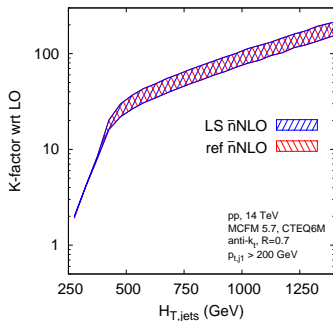
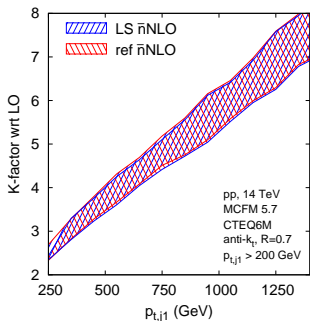
Reference-observable method

Take a reference observable identical at LO to the observable A

$$\begin{aligned}\sigma_{Z+j@NNLO}^{(A)} &= \sigma_{Z+j@NNLO}^{(\text{ref})} + (\sigma^{(A)} - \sigma^{(\text{ref})})_{Z+j@NNLO} \\ &= \sigma_{Z+j@NNLO}^{(\text{ref})} + (\sigma^{(A)} - \sigma^{(\text{ref})})_{Z+2j@NLO}\end{aligned}$$

If the reference observable converges well

$$\sigma_{Z+j@NNLO}^{(A)} \simeq \sigma_{Z+j@NLO}^{(\text{ref})} + (\sigma^{(A)} - \sigma^{(\text{ref})})_{Z+2j@NLO}$$

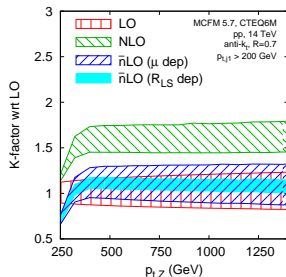


Z+jet at NLO

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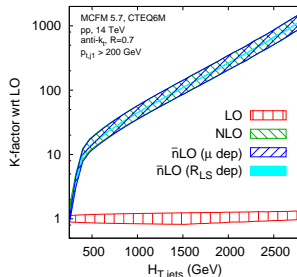
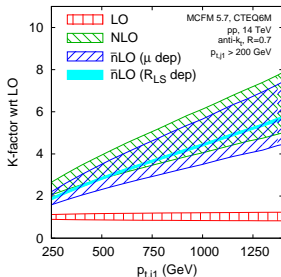
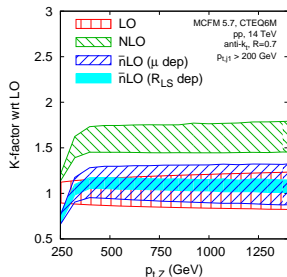
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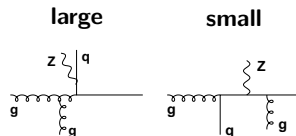
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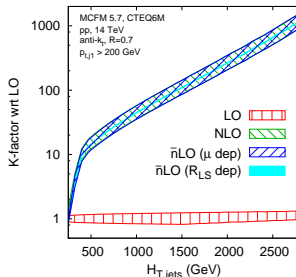
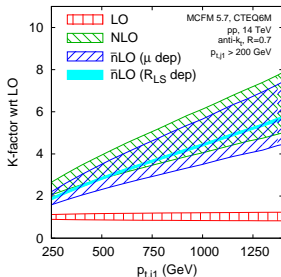
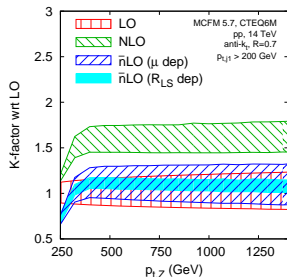


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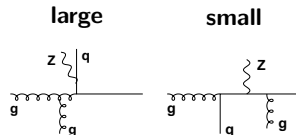
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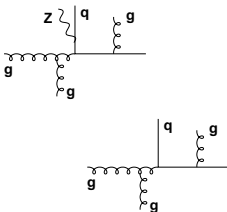
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- small R uncertainties – driven only by subleading diagrams



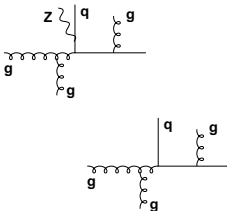
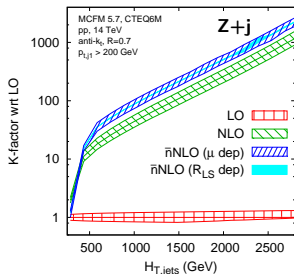
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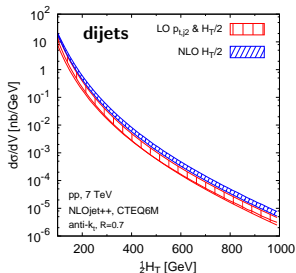
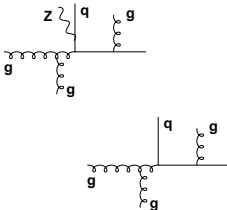
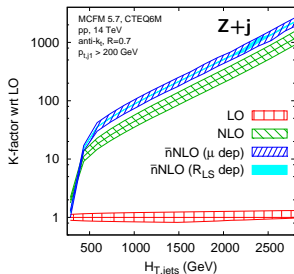
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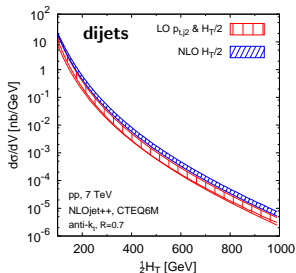
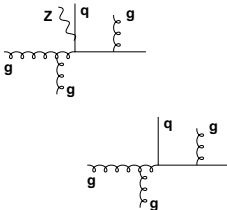
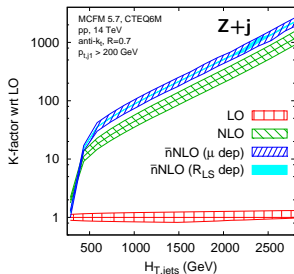
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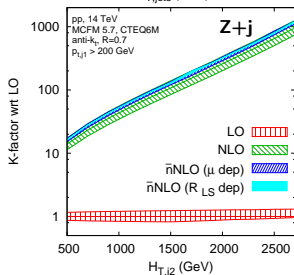
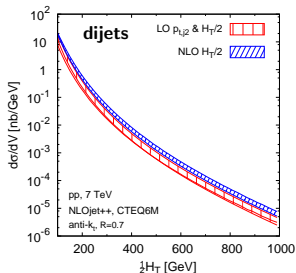
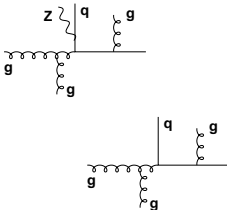
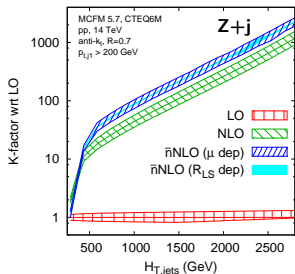
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 - ▶ and indeed it is small!